

# Superatom-assembled boranes, carboranes, and low-dimensional boron nanomaterials based on aromatic icosahedral B<sub>12</sub> and C<sub>2</sub>B<sub>10</sub>

Qiao-Qiao Yan, Yan-Fang Wei, Qiang Chen (X), Yue-Wen Mu (X), and Si-Dian Li (X)

Institute of Molecular Science, Shanxi University, Taiyuan 030006, China

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# ABSTRACT

Using the experimentally known aromatic icosahedral superatoms  $I_h B_{12}H_{12}^{2-}$  and  $D_{5d} 1,12$ - $C_2B_{10}H_{12}$  as building blocks and based on extensive density functional theory calculations, we predict herein a series of core–shell superpolyhedral boranes and carboranes in a bottom-up approach, including the high-symmetry  $T_h B_{12}@B_{152}H_{72}^{2-}$  (2),  $C_{2h} C_2B_{10}@B_{152}H_{72}$  (3),  $D_{3d} B_{12}@B_{144}H_{66}$ (4),  $I_h B_{12}@C_{24}B_{120}H_{72}^{2-}$  (6), and  $D_{5d} C_2B_{10}@C_{24}B_{120}H_{72}$  (7). More interestingly, the superatom-assembled linear  $D_{2h} B_{36}H_{32}^{2-}$  (8), close-packed planar  $D_{3d} B_{84}H_{60}^{2-}$  (10), and nearly close-packed core–shell  $D_{3d} B_{12}@B_{144}H_{66}$  (4) can be extended periodically to form the one-dimensional (1D)  $\alpha$ -rhombohedral borane nanowire  $B_{12}H_{10}$  (*Pmmm*) (9), two-dimensional (2D)  $\alpha$ -rhombohedral monolayer borophane  $B_{12}H_6$  ( $P\overline{3}m1$ ) (11), and the experimentally known three-dimensional (3D)  $\alpha$ -rhombohedral boron ( $R\overline{3}m$ ) (12) which can be viewed as an assembly of the monolayer  $B_{12}H_6$  (11) staggered in vertical direction, setting up a bottom-up strategy to form low-dimensional boron-based nanomaterials from their borane "seeds" via partial or complete dehydrogenations. Detailed bonding analyses indicate that the high stability of these nanostructures originates from the spherical aromaticity of their icosahedral  $B_{12}$  or  $C_2B_{10}$  structural units which possess the universal skeleton electronic configuration of  $1S^21P^e1D^{10}1F^e$  following the Wade's n+1 rule. The infrared (IR) and Raman spectra of the most-concerned neutral  $B_{12}@B_{144}H_{66}$  (4) and  $C_2B_{10}@C_{24}B_{120}H_{72}$ (7) are computationally simulated to facilitate their experimental characterizations.

# **KEYWORDS**

superatom-assembling, bottom-up approach, density functional theory, superpolyhedral boranes, superpolyhedral carboranes, borophanes

# 1 Introduction

Increasing experimental and theoretical evidences in the past two decades suggest that cluster superatoms may serve as building blocks to form self-assembled compounds and nanomaterials in bottom-up approaches [1-4]. The best experimentally known superatoms in boron chemistry are the 12-vertex icosahedral borane  $I_h B_{12}H_{12}^{2-}$  (<sup>1</sup>A<sub>g</sub>) and icosahedral carborane  $D_{5d}$  1,12- $C_2B_{10}H_{12}$  (<sup>1</sup>A<sub>1g</sub>) due to their superb chemical stabilities [5–10]. Both these icosahedral species prove to follow the superatomic configuration of 1S21P61D101F8 (see Fig. S1 in the Electronic Supplementary Material (ESM)) and match the Wade's n+1 rule in skeleton-electron counts [11], making them spherically aromatic in nature. It is also well known that all the seventeen bulk boron allotropes experimentally observed to date contain icosahedral B12 structural units which, in most cases, are accompanied by certain numbers of boron atoms lying outside the cages as interstitial atoms [12, 13], except the simplest form, arhombohedral boron (a-B12) firstly synthesized in 1958 [14], which contains a D<sub>3d</sub> B<sub>12</sub> icosahedral unit cell in a nearly closepacked arrangement. In fact, with six B-B 2c-2e  $\sigma$  bonds on the top and bottom and six B–B–B 3c–2e  $\sigma$  bonds on the waist, each  $D_{3d}$  B<sub>12</sub> unit cell in  $\alpha$ -B<sub>12</sub> matches the Wade's *n*+1 rule as a spherically aromatic structural unit as discussed in details late, making  $\alpha$ -B<sub>12</sub> unique in boron bulk allotropes. Monolayer  $\gamma$ -B<sub>28</sub> composed of icosahedral B12 clusters and B2 dumbbells was synthesized in 2015 on copper foils by chemical vapor deposition at low pressures with hydrogen as carrier gases [15]. However, small bare  $B_n^{-10}$  clusters (n < 19) characterized in gas phases by persistent joint photoelectron spectroscopy (PES) and firstprinciples theory investigations in the past two decades all possess planar or quasi- planar structures, with their cage-like isomers with unsaturated dangling bands on the surfaces being much less stable in thermodynamics [16–22]. Neutral  $B_{12}$  as the bare boron cluster most-concerned in this work has been experimentally confirmed to have a quasi-planar  $C_{3\nu}$  structure analogous to benzene ( $C_6H_6$ ) in  $\pi$  bonding [16, 20]. The smallest bare cage-like  $B_n$  cluster characterized in PES experiments is the seashell-like  $C_2$  $B_{28}^{0/-}$  [23], followed by chiral cage-like  $C_3/C_2 B_{39}^{-}$  [24], cage-like  $D_{2d}$  $B_{40}^{-10}$  [25], and bilayer  $D_{2h} B_{48}^{-10}$  [26]. Obviously, how to form bulk boron allotropes from bare  $B_n$  clusters or closo-icosahedral boranes in a bottom-up approach still remains unknown to date in both experiments and theory.

Keeping the inspiration in mind and based on icosahedral  $B_{12}H_{12}^{2-}$  and  $1,12-C_2B_{10}H_{12}$ , we predict in this work at densityfunctional theory (DFT) level a series of superatom-assembled core-shell superpolyhedral boranes and carboranes  $T_{h}$  $B_{12}@B_{152}H_{72}^{2-}$  (2),  $C_{2h} C_2B_{10}@B_{152}H_{72}$  (3),  $D_{3d} B_{12}@B_{144}H_{66}$  (4),  $I_h$  $B_{12}@C_24B_{120}H_{72}^{2-}$  (6), and  $D_{5d} C_2B_{10}@C_{24}B_{120}H_{72}$  (7) in the structural motif of  $C_{60}$ . The linear  $D_{2h} B_{36}H_{32}^{2-}$  (8), close-packed planar  $D_{3d} B_{84}H_{60}^{2-}$  (10), and core-shell  $D_{3d} B_{12}@B_{144}H_{66}$  (4)



assembled in similar patterns have been extended to form the infinite one-dimensional (1D)  $\alpha$ -rhombohedral borane nanowire  $B_{12}H_{10}$  (*Pmmm*) (9), two-dimensional (2D)  $\alpha$ -rhombohedral monolayer borophane  $B_{12}H_6$  ( $P\overline{3}m1$ ) (11), and three-dimensional (3D)  $\alpha$ -rhombohedral boron ( $R\overline{3}m$ ) ( $\alpha$ - $B_{12}$ ) (12), presenting a bottom-up approach to form borane and carborane nanoclusters and their corresponding low-dimensional nanomaterials by superatom-assembling via partial or complete dehydrogenations.

#### 2 Theoretical procedures

The manually constructed superatom-assembled superpolyhedral boranes and carboranes in the structural motifs of C60 via partialdehydrogenizations were fully optimized at both the hybrid Perdew-Burke-Ernzerhof (PBE0)/6-31G(d) and B3LYP/6-31G(d) levels [27, 28] in Fig. 1, with both their vibrational frequencies and stabilities of wavefunctions checked to ensure that all obtained structures are true minima of the systems. The high-symmetry core-shell 2, 3, 4, 6, and 7 in Fig.1 and linear  $B_{36}H_{32}$ <sup>2-</sup> (8) and close-packed planar  $B_{84}H_{60}^{2-}$  (10) in Fig. 2 appeared to be much more stable (at least 2.10 eV) than their distorted triplet and quintet counterparts at PBE0. As a demonstration, the perfect superpolyhedral borane  $I_h C_{24}B_{120}H_{84}$  (5) (<sup>1</sup>A<sub>g</sub>) [29] was manually constructed by substituting the twelve C5 pentagons on the surface of  $C_{60}$  with twelve  $D_{5d}$  1,12- $C_2B_{10}H_{12}$  superatoms with the  $C_5$ molecular axis located exactly perpendicular to the cage surface towards the cage center, leading to the formation of thirty intercage B–B  $\sigma$  bonds and release of thirty H<sub>2</sub> molecules. The perfect core-shell  $I_h$  B<sub>12</sub>@C<sub>24</sub>B<sub>120</sub>H<sub>72</sub><sup>2-</sup> (6) (<sup>1</sup>A<sub>g</sub>) with an icosahedral B<sub>12</sub> core covalently saturated by twelve inter-cage B-C  $\sigma$  bonds (~ 1.60 Å) was obtained by incorporating an  $I_h B_{12}H_{12}^{2-}$  at the center of  $C_{24}B_{120}H_{84}$  (5) to release 6 H<sub>2</sub> molecules, while the core-shell  $D_{5d}$  $C_2B_{10}@C_{24}B_{120}H_{72}$  (7) (<sup>1</sup>A<sub>o</sub>) with an icosahedral  $D_{5d}$   $C_2B_{10}$  core was achieved by incorporating a  $D_{5d}$  C<sub>2</sub>B<sub>10</sub> core at the center of the C24B120H84 (5) shell. Nucleus-independent chemical shifts (NICS) [30, 31] were calculated at the geometrical centers to check the spherical aromaticity of systems, with the corresponding NICS-ZZ distributions obtained using the Multiwfn 3.8 code [32]. The infrared (IR) and Raman spectra of neutral B12@B144H66 (4) and  $C_2B_{10}@C_{24}B_{120}H_{72}$  (7) were simulated at PBE0/6-31G(d).

Structural optimizations on 1D, 2D, and 3D crystals were carried out using the Vienna *ab initio* simulation package (VASP 5.4) [33, 34] at DFT level within the PBE parametrization of generalized gradient approximation (GGA) [35]. A vacuum slab

(more than 15 Å) was selected to prevent mirror interactions and a plane-wave basis set with an energy cutoff of 500 eV was used. The Brillouin zones were sampled with  $2\pi \times 0.02$  Å<sup>-1</sup> spacing in the reciprocal space by the Monkhorst–Pack scheme [36]. Atomic structures were fully relaxed using the conjugate gradient method until the maximum force on each atom was less than  $10^{-3}$  eV·Å<sup>-1</sup> and the energy precision was set to  $10^{-5}$  eV per atom. The phonon dispersion spectra of 1D B<sub>12</sub>H<sub>10</sub> (9), 2D B<sub>12</sub>H<sub>6</sub> (11), and 3D α-B<sub>12</sub> (12) were calculated using the density-functional perturbation theory (DFPT) method implemented in the Phonopy program [37] interfaced with VASP, with their band structures obtained at the hybrid HSE06 level [38] (Fig. 3).

Born-oppenheimer molecular dynamics (BOMD) simulations were performed on the concerned species for 30 ps using the CP2K program [39], with the Goedecker–Teter–Hutter (GTH)-PBE pseudopotential and DZVP-MOLOPT-SR-GTH basis sets used for B, C, and H. Chemical bonding analyses were performed on isolated clusters using the adaptive natural density partitioning (AdNDP) approach [16, 40] and on 1D, 2D, and 3D crystals utilizing the solid-state AdNDP (SSAdNDP) [41]. The results were visualized using the VMD 1.9.3 program [42]. The calculations on zero-dimensional clusters were performed utilizing the Gaussian 16 software [43] and on low-dimensional crystals employing the VASP software.

# 3 Results and discussions

#### 3.1 Structures and stability

The superatom-assembled high-symmetry superpolyhedral borane  $T_h B_{152}H_{84}$  (1) and carborane  $I_h C_{24}B_{120}H_{84}$  (5) and their core-shell superpolyhedral counterparts  $T_h B_{12} @B_{152} H_{72}^{2-}$  (2),  $C_{2h}$  $C_2B_{10}@B_{152}H_{72}$  (3),  $D_{3d} B_{12}@B_{144}H_{66}$  (4),  $I_h B_{12}@C_{24}B_{120}H_{72}^{2-}$  (6), and  $D_{5d} C_2 B_{10} @C_{24} B_{120} H_{72}$  (7) optimized at both PBE0/6-31G(d) and B3LYP/6-31G(d) levels are collectively depicted in Fig.1. With the valences of all the component surface atoms fully satisfied by either terminal -H atoms, bridging -H- atoms, or intercage B–B or B–C  $\sigma$  interactions on the cage surfaces, the obtained 1-7 with stable wavefunctions all turn out to be true minima of the systems without imaginary vibrational frequencies, as tabulated in Table 1. The suprapolyhedral boranes and carboranes with a closo-B<sub>12</sub> or closo-C<sub>2</sub>B<sub>10</sub> core at the center possess the large highest occupied molecular orbital (HOMO)-lowest unoccupied molecular orbital (LUMO) gaps between  $E_{gap} = 2.85-7.09$  eV



**Figure 1** Optimized structurers of the superatom-assembled  $T_h B_{152}H_{84}$  (1) and  $I_h C_{24}B_{120}H_{84}$  (5) and their core-shell superpolyhedral counterparts  $T_h B_{12}@B_{152}H_{72}^{2-}$  (2),  $C_{2h} C_{2B_{10}}@B_{152}H_{72}$  (3),  $D_{3d} B_{12}@B_{144}H_{66}$  (4),  $I_h B_{12}@C_{24}B_{120}H_{72}^{2-}$  (6), and  $D_{5d} C_{2B_{10}}@C_{24}B_{120}H_{72}$  (7) at PBE0/6-31G(d) and B3LYP/6-31G(d) levels using the icosahedral  $B_{12}H_{12}^{2-}$  and 1,12- $C_{2B_{10}}H_{12}$  as building blocks in the structural motif of  $I_h C_{60}$  via partial dehydrogenations (pink B, black C, and white H, with the 12-vertex polyhedrons highlighted in purple).



**Figure 2** Optimized superatom-assembled linear  $D_{2h} B_{36} H_{32}^{2-}$  (8), close-packed hexagonal  $D_{3d} B_{84} H_{60}^{2-}$  (10), and nearly close-packed core-shell  $D_{3d} B_{12} @B_{144} H_{66}$  (4) and their extended periodic nanomaterials 1D α-rhombohedral borane nanowire  $B_{12} H_{10}$  (*Pmmm*) (9), 2D α-rhombohedral monolayer borophane  $B_{12} H_6$  ( $P\bar{3}m1$ ) (11), and 3D nearly close-packed α-rhombohedral boron ( $R\bar{3}m$ ) (12) at GGA-PBE level, with the optimized lattice parameters indicated.

Table 1 Calculated lowest vibrational frequencies  $v_{min}$  (cm<sup>-1</sup>), HOMO–LUMO energy gaps  $\Delta E_{gap}$  (eV), and NICS (ppm) values at the geometrical centers of 1–8 and 10

| Cluster  | Symmetry | $v_{\min}$ (cm <sup>-1</sup> ) | $\Delta E_{\rm gap}$ (eV) | NICS (ppm) |
|--|----------|--------------------------------|---------------------------|------------|
| $B_{152}H_{84}(1)$                                       | $T_h$    | 133.5                          | 4.37                      | -2.23      |
| $B_{12}@B_{152}H_{72}^{2-}$ (2)                          | $T_h$    | 82.6                           | 3.58                      | -17.79     |
| $C_{2}B_{10}@B_{152}H_{72}$ (3)                          | $C_{2h}$ | 132.3                          | 3.83                      | -20.25     |
| $B_{12}@B_{144}H_{66}(4)$                                | $D_{3d}$ | 86.21                          | 2.85                      | -19.92     |
| $C_{24}B_{120}H_{84}$ (5)                                | $I_h$    | 127.0                          | 6.59                      | -3.51      |
| $B_{12} @ C_{24} B_{120} H_{72}{}^{2-} (\boldsymbol{6})$ | $I_h$    | 140.3                          | 7.09                      | -27.22     |
| $C_{2}B_{10}@C_{24}B_{120}H_{72}(7)$                     | $D_{5d}$ | 141.7                          | 6.04                      | -27.94     |
| B <sub>36</sub> H <sub>32</sub> <sup>2-</sup> (8)        | $D_{2h}$ | 50.50                          | 6.07                      | -27.99     |
| $B_{84}H_{60} \sim (10)$                                 | $D_{3d}$ | 67.53                          | 4.91                      | -22.17     |

(which are even higher than the corresponding value (3.04 eV) of  $I_h C_{60}$  at the same theoretical level, except for  $B_{12}@B_{144}H_{66}$  (4)) and negative calculated NICS values between NICS = -17.79--35.30 ppm at the geometrical centers (Table 1), indicating that they are all spherically aromatic in nature and highly chemically stable.

It is also noticed that the core–shell superpolyhedral boranes and carboranes (**2**, **3**, **4**, **6**, and **7**) have significantly larger negative NICS values than their corresponding empty shells (1 and 5), indicating that both the closo– $B_{12}$  and closo– $C_2B_{10}$  cores help to improve the overall spherical aromaticity of the systems effectively.

The neutral  $T_h B_{152}H_{84}$  (1) in Fig. 1 is assembled from twelve  $I_h B_{12}H_{12}^{2-}$  icosahedra in the structural motif of  $C_{60}$  by releasing thirty  $H_2$  molecules, with eight hexacoordinate B atoms in eight  $\eta^6$ -B<sub>6</sub> hexagons added on the cage surface to compensate for the 24 valence electrons deficiency of the system according to Wade's n+1 rule. The 1+12 core–shell suprapolyhedral borane  $T_h B_{12}@B_{152}H_{72}^{2-}$  (2) is constructed by insertion of an icosahedral  $B_{12}$  core at the cage center of  $B_{152}H_{84}$  (1) which is covalently satisfied by twelve inter-cage B–B  $\sigma$  single bonds (~ 1.63 Å) in radial directions with the twelve  $B_{12}$  icosahedral units evenly distributed on the cage surface in the first coordination sphere. The  $T_h B_{12}^{2-}$ 

core at the center of the dianion (2) follows the Wade's n+1 rule with two extra electrons. Using a para-icosahedral  $1,2-C_2B_{10}$  to replace the  $B_{12}$  core, the neutral 1+12 suprapolyhedral carborane  $C_{2h}C_2B_{10}@B_{152}H_{72}$  (3) is generated.

Eliminating the eight hexacoordinate B atoms on the cage surface of  $B_{12}@B_{152}H_{72}^{2-}$  (2) and six H atoms from the top and bottom while assigning six B–H–B 3c–2e bonds in a  $D_{3d}$ symmetry between neighboring B12 icosahedra on the waist generate the nearly close-packed core-shell D<sub>3d</sub> B<sub>12</sub>@B<sub>144</sub>H<sub>66</sub> (4) which, as a true minimum of the neutral, contains a  $D_{3d}$  B<sub>12</sub> core with six equivalent B–B 2c–2e  $\sigma$  bonds ( $r_{B-B} = 1.69$  Å) in vertical direction and six equivalent B-B-B 3c-2e bonds with six B<sub>12</sub> icosahedrons around it on the waist. Encouragingly, the thusobtained  $B_{12}@B_{144}H_{66}$  (4) contains a  $D_{3d}B_{12}$  core which has exactly the same symmetry as the  $D_{3d}$  B<sub>12</sub> unit cell in the experimentally known  $\alpha$ -B<sub>12</sub> (12) and possesses the same coordination environment with the twelve B<sub>12</sub> icosahedra in the first coordination sphere as that of a  $D_{3d}$  B<sub>12</sub> unit cell in  $\alpha$ -B<sub>12</sub> boron (12).  $B_{12}@B_{144}H_{66}$  (4) can thus be viewed as the core-shell suprapolyhedral borane "seed" of 3D  $\alpha$ -B<sub>12</sub> (12) in a bottom-up approach by a periodic expansion in three dimensions via a complete dehydrogenation, as demonstrated in Fig. 2.  $B_{12}@B_{144}H_{66}$  (4) is the first suprapolyhedral borane cluster reported to date which can serve as the "seed" of an experimentally known bulk boron allotrope in a bottom-up strategy.

As mentioned above,  $D_{5d}$  1,12- $C_2B_{10}H_{12}$  can be used as building blocks to form the perfect cage-like superpolyhedral carborane  $I_h$  $C_{24}B_{120}H_{84}$  (5) (<sup>1</sup>A<sub>g</sub>) which can be expanded to the 1+12 core-shell superpolyhedral carborane  $I_h$   $B_{12}@C_{24}B_{120}H_{72}^{-2}$  (6) (<sup>1</sup>A<sub>g</sub>) with an icosahedral  $I_h$   $B_{12}$  core at the center and  $D_{5d}$   $C_2B_{10}@C_{24}B_{120}H_{72}$  (7) (<sup>1</sup>A<sub>g</sub>) with an icosahedral  $D_{5d}$   $C_2B_{10}$  core at the center, respectively. As shown in Table 1, with  $v_{min} = 140.3$ , 141.7 cm<sup>-1</sup>,  $\Delta E_{gap} = 7.09$ , 6.04 eV, and NICS = -27.22, -27.94 ppm, respectively, both the 1+12 core-shell  $B_{12}@C_{24}B_{120}H_{72}^{-2}$  (6) and  $C_2B_{10}@C_{24}B_{120}H_{72}$  (7) as true minima of the systems are spherically aromatic and highly chemically stable. It is easy to understand that both the twelve  $C_2B_{10}$  icosahedra on the cage surface and the  $C_2B_{10}$  icosahedral core at the center in  $C_2B_{10}@C_{24}B_{120}H_{72}$  (7) match the Wade's n+1rule, while with two extra electrons, the icosahedral  $B_{12}^{2-}$  core in the dianionic  $B_{12}@C_{24}B_{120}H_{72}^{-2}$  (6) also behaves like a superatom.

Following a similar superatom-assembling pattern, we designed the linear  $D_{2h} B_{36}H_{32} \sim (8)$  in Fig. 2 which contains two B–H–B 3c–2e bonds between neighboring B<sub>12</sub> units. The central B<sub>12</sub>H<sub>10</sub> unit with four B–H–B 3c–2e bonds in 8 matches the Wade's *n*+1 rule due to the fact that the formation of one B–H–B 3c–2e bond contributes half a skeleton electron to the related B<sub>12</sub> unit. It turns out to be true, as shown in the bonding analyses detailed in Fig. 4. Intriguingly, as shown in Fig. 2, the linear  $B_{36}H_{32}^{2-}$  (8) can be extended periodically along the  $C_2$  molecular axis to form the infinite 1D  $\alpha$ -rhombohedral borane nanowire  $B_{12}H_{10}$  (*Pmmm*) (9) with the calculated lattice parameter of a = 4.90 Å at GGA-PBE in which each  $B_{12}H_{10}$  unit cell forms four B–H–B 3c–2e bonds with its two equivalent  $B_{12}H_{10}$  neighbors. 1D  $B_{12}H_{10}$  (9) as a true minimum of the system without imaginary phonon frequencies (Fig. 3(a)) can thus be viewed as a superatom-assembled monochain borane nanowire. It turns out to be a semiconductor in nature with the direct band gap of 5.54 eV at HSE06, as shown in its calculated band structures and projected densities of states (PDOS) depicted in Fig. 3.

The superatom-assembled close-packed perfect hexagonal  $D_{3d}$ B<sub>84</sub>H<sub>60</sub><sup>2-</sup> (**10**) contains a  $D_{3d}$  B<sub>12</sub>H<sub>6</sub> icosahedron at the center surrounded by six equivalent periphery B<sub>12</sub>H<sub>9</sub> units around it. As highlighted in blue in Fig. 2, the central B<sub>12</sub>H<sub>6</sub> unit forms six equivalent B–B–B 3c–2e bonds with its six periphery B<sub>12</sub> icosahedral ligands on the mirror plane in a  $D_{3d}$  distribution. Such a hexagonal planar dianion can be extended in two dimensions to form the perfect 2D  $\alpha$ -rhombohedral monolayer borophane B<sub>12</sub>H<sub>6</sub> ( $P\overline{3}m1$ ) (**11**) without imaginary phonon frequencies (Fig. 3(b)) which, as a semiconductor with the calculated indirect bond gap of 2.25 eV at HSE06, is a close-packed 2D self-assembled monolayer (SAM) borophane with the optimized lattice parameters of a = b = 4.87 Å, and  $\gamma = 120^{\circ}$  at GGA-PBE (Fig. 3). We also notice that the  $D_{3d}$  B<sub>12</sub>H<sub>6</sub> core in the close-packed B<sub>84</sub>H<sub>60</sub><sup>2-</sup>



Figure 3 Calculated band structures and PDOS, and phonon dispersion spectra of (a) 1D  $B_{12}H_{10}$  (9), (b) 2D  $B_{12}H_6$  (11), and (c) 3D  $\alpha$ - $B_{12}$  (12) at the hybrid HSE06 level.

(10) can be removed from the dianion to form a perfect hexagonal ring  $D_{3d} B_{72}H_{60}$  (13) which is a true minimum of the system analogous to benzene (Fig. S2 in the ESM).

More profoundly, the core-shell  $D_{3d}$   $B_{12}@B_{144}H_{66}$  (4) can be expanded in three dimensions in a bottom-up approach via a complete dehydrogenation to form the experimentally known nearly close-packed 3D  $\alpha$ -B<sub>12</sub> ( $R\overline{3}m$ ) (12) without imaginary phonon frequencies (Fig. 3(c)) which can also be viewed as an assembly of the monolayer  $B_{12}H_6$  (11) staggered in vertical direction. With the optimized lattice parameters of a = b = c =5.01 Å and  $\alpha = \beta = \gamma = 57.99^{\circ}$  at GGA-PBE (Fig. 2) and indirect band gap of 1.85 eV at HSE06 (close to the experimental value of 2.0 eV [44]) (Fig. 3), 3D  $\alpha$ -B<sub>12</sub> (12) has the calculated B–B bond lengths of  $r_{\rm B-B} = 1.66$  Å between the B<sub>12</sub> unit cell and its six B<sub>12</sub> icosahedral neighbors on the vertical direction and the B-B distances of  $r_{B-B} = 1.98$  Å with its six  $B_{12}$  icosahedral  $B_{12}$  neighbors in the horizontal plane where each  $D_{3d}$  B<sub>12</sub> unit cell forms six equivalent B–B–B 3c–2e  $\sigma$ -bonds with its six  $D_{3d}$  B<sub>12</sub> neighbors. Thus, as outlined in Fig. 2, starting from a perfect icosahedral  $I_h$  $B_{12}H_{12}^{2-}$ , the 1D mono-chain nanowire  $B_{12}H_{10}$  (9), 2D SAM  $B_{12}H_{6}$ (11), and 3D  $\alpha$ -B<sub>12</sub> (12) can be systematically assembled in a bottom-up approach via partial or complete dehydrogenations.

Extensive Born–Oppenheimer molecular dynamics (BOMD) simulations on  $B_{12}@B_{152}H_{72}^{2-}$  (2) at 500K,  $B_{12}@B_{144}H_{66}$  (4) at 1000 K, and  $C_2B_{10}@C_{24}B_{120}H_{72}$  (7) at 1200 K in Fig.S3 in the ESM clearly indicate that these core–shell suprapolyhedral species are all highly dynamically stable, as evidenced by their small calculated average root-mean-square-deviations of RMSD = 0.08/0.10/0.09 Å and maximum bond length deviations of MAXD = 0.63/0.39/0.40 Å, respectively. These core–shell structures are all well maintained during BOMD simulations in 30 ps, with no low-lying isomers observed. We also assess the dynamic stability of 1D  $B_{12}H_{10}$  (9) at

800 K and 2D  $B_{12}H_6$  (11) at 1000 K by performing BOMD simulations on a 1 × 1 × 4 and 4 × 4 × 1 supercell for 30 ps (Fig. S3 in the ESM). With the calculated RMSD = 0.09/0.10 Å and MAXD = 0.31/0.39 Å, respectively, these semiconducting crystals all appear to be highly dynamically stable at high temperatures.

## 3.2 Bonding pattern analyses

Detailed AdNDP and SSAdNDP bonding analyses unveil both the localized and delocalized bonds in the concerned zerodimensional (0D), 1D, 2D, and 3D species [40, 41] and help to comprehend the bottom-up approach employed in this work. The spherically aromatic icosahedral Ih B12H12<sup>2-</sup> and D5d C2B10H12 with the superatomic electronic configuration of  $1S^2 1P^6 1D^{10} 1F^8$  are wellknown to follow the Wade's n+1 rule (Fig. S1 in the ESM). Interestingly, as can be seen from Figs. 4(a) and 4(d), both the central  $D_{2h}$  B<sub>12</sub>H<sub>10</sub> unit in B<sub>36</sub>H<sub>32</sub><sup>2-</sup> (8) and  $D_{2h}$  B<sub>12</sub>H<sub>10</sub> unit cell in 1D nanowire B<sub>12</sub>H<sub>10</sub> (9) possess 1 S-type bond, 3 P-type bonds, 5 D-type bonds, and 4 F-type bonds following the superatomic skeleton electronic configuration of 1S21P61D101F8. Similarly, the central  $D_{3d}$  B<sub>12</sub>H<sub>6</sub> unit in B<sub>84</sub>H<sub>60</sub><sup>2-</sup> (10) and  $D_{3d}$  B<sub>12</sub>H<sub>6</sub> unit cell in 2D borophane  $B_{12}H_6$  (11) contain 1 S-type bond, 3 P-type bonds, 5 D-type bonds, and 4 F-type skeleton bonds too, again matching the superatomic bonding pattern (Figs. 4(b) and 4(e)). More intriguingly, as demonstrated in Figs. 4(c) and 4(f), in addition to the 6 equivalent B-B 2c-2e  $\sigma$  bonds with six B<sub>12</sub> icosahedra in vertical direction and 6 equivalent B-B-B 3c-2e bonds with six  $B_{12}$  icosahedra on the waist, both the central  $D_{3d}$   $B_{12}$  unit in  $B_{12}@B_{144}H_{66}$  (4) and  $D_{3d}B_{12}$  unit cell in 3D  $\alpha$ - $B_{12}$  (12) also possess 1 S-type bond, 3 P-type bonds, 5 D-type bonds, and 4 F-type skeleton bonds, again matching the Wade's rule. The 0D, 1D, 2D, and 3D species in Figs. 1 and 2 can thus all be viewed as lowdimensional nanostructures assembled in bottom-up approaches



**Figure 4** Comparison of the AdNDP bonding patterns of (a)  $D_{2h}B_{36}H_{32}^{2-}(\mathbf{8})$ , (b)  $D_{3d}B_{84}H_{60}^{2-}(\mathbf{10})$ , (c)  $D_{3d}B_{124}B_{66}(\mathbf{4})$ , (d) 1D  $B_{12}H_{10}(\mathbf{9})$ , (e) 2D  $B_{12}H_6(\mathbf{11})$ , and (f) 3D  $\alpha$ -rhombohedral boron (12). For brevity, only the chemical bonds related to the central  $B_{12}$  units in  $B_{36}H_{32}^{2-}(\mathbf{8})$ ,  $B_{84}H_{60}^{2-}(\mathbf{10})$ , and  $B_{12}@B_{144}H_{66}(\mathbf{4})$  are depicted here, with their complete bonding patterns presented in Fig. S4 in the ESM.

The calculated NICS-ZZ distributions depicted in Fig. 5 further evidence the superatomic aromatic nature of the icosahedral  $B_{12}$ and  $C_2B_{10}$  units in  $D_{2h}$   $B_{36}H_{32}^{-2-}$  (**8**),  $D_{3d}$   $B_{84}H_{60}^{2-}$  (**10**),  $D_{3d}$  $B_{12}@B_{144}H_{66}$  (**4**), and  $D_{5d}$   $C_2B_{10}@C_{24}B_{120}H_{72}$  (**7**) where the NICS-ZZ value at the center of each  $B_{12}$  or  $C_2B_{10}$  icosahedron is highly negative. Similar negative NICS-ZZ distributions are expected to exist in 1D  $B_{12}H_{10}$  (**9**), 2D  $B_{12}H_6$  (**11**), and 3D  $\alpha$ - $B_{12}$  (**12**) where each icosahedral  $B_{12}$  unit cell follows the skeleton electronic configuration of  $1S^{2}1P^{6}1D^{10}1F^{8}$  and matches the Wade's n+1 rule, as demonstrated in details in Figs. 4(d)–4(f), respectively.

#### 3.3 Spectral simulations

The IR and Raman spectra of the most-concerned neutral core–shell superpolyhedral borane and carborane were computationally simulated in Fig. 6 to facilitate their future experimental characterizations.  $D_{3d}$  B<sub>12</sub>@B<sub>144</sub>H<sub>66</sub> (4) exhibits two strong IR absorptions at 1890 cm<sup>-1</sup> (a<sub>g</sub>) and 2706 cm<sup>--1</sup> (a<sub>g</sub>) and three Raman active peaks at 211 cm<sup>-1</sup> (a<sub>g</sub>), 2707 cm<sup>-1</sup> (a<sub>g</sub>), and 2765 cm<sup>-1</sup> (a<sub>g</sub>), with a 183 cm<sup>-1</sup> (a<sub>g</sub>) symmetrical vibration corresponding to typical radial breathing modes (RBMs) of the system. The major IR active peak of  $D_{5d}$  C<sub>2</sub>B<sub>10</sub>@C<sub>24</sub>B<sub>120</sub>H<sub>72</sub> (7) occurs at 2735 cm<sup>-1</sup> (a<sub>10</sub>). The four vibration modes at 207 cm<sup>-1</sup>



(c)  $D_{3d} B_{12} @ B_{144} H_{66}$  (4) (d)  $D_{5d} C_2 B_{10} @ C_{24} B_{120} H_{72}$  (7) **Figure 5** NICS-ZZ distributions of (a)  $D_{2h} B_{36} H_{32^2}$  (8), (b)  $D_{3d} B_{84} H_{60^2}$  (10), (c)  $D_{3d} B_{12} @ B_{144} H_{66}$  (4), and (d)  $D_{5d} C_2 B_{10} @ C_{24} B_{120} H_{72}$  (7) at PBE0/6-31G(d) level.

 $(a_{1g})$ , 253 cm<sup>-1</sup>  $(a_{1g})$ , 2774 cm<sup>-1</sup>  $(a_{1g})$ , and 3276 cm<sup>-1</sup>  $(a_{1g})$  could be used as Raman fingerprints to characterize the  $D_{5d}$  species. The strong Raman peaks 253 cm<sup>-1</sup>  $(a_{p})$  represents its typical RBM.



Figure 6 Simulated IR and Raman spectra of (a)  $D_{3d}B_{12}@B_{144}H_{66}$  (4) and (b)  $D_{5d}C_2B_{10}@C_{24}B_{120}H_{72}$  (7) at PBE0/6-31G(d) level.

### 4 Conclusions

In summary, a series of superatom-assembled boranes and carboranes  $T_h B_{12} @B_{152} H_{72}^{2-}$  (2),  $C_{2h} C_2 B_{10} @B_{152} H_{72}$  (3),  $D_{3d} B_{12} @$  $B_{144}H_{66}$  (4),  $I_h B_{12}@C_{24}B_{120}H_{72}^{2-}$  (6),  $D_{5d} C_2B_{10}@C_{24}B_{120}H_{72}$  (7),  $D_{2h}$  $B_{36}H_{32}^{2-}$  (8), and  $D_{3d}$   $B_{84}H_{60}^{2-}$  (10) have been predicted at DFT in this work which can be extended to form the 1D borane nanowire  $B_{12}H_{10}$  (9), 2D monolayer borophane  $B_{12}H_6$  (11), and 3D  $\alpha$ - $B_{12}$ (12) boron semiconductors via partial or complete dehydrogenizations. Detailed bonding analyses indicate that the high stability of these superatom-assembled nanostructures can be traced back to the spherical aromaticity of their icosahedral B<sub>12</sub> or  $C_2B_{10}$  structural units which follow the Wade's n+1 rule in skeleton electron counts. Similar to graphene, the monolayer  $B_{12}H_6$  (11) could be rolled into 1D  $\alpha$ -rhombohedral borane nanotubes or assembled into 2D few-layered a-rhombohedral borophanes staggered in vertical direction. Such superatomassembled nano-structures based on icosahedral closo-B<sub>12</sub> and  $closo-C_2B_{10}$  are expected to be realized in future experiments to extend the structural domain of boron-based nanomaterials.

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